

## 發電機原理

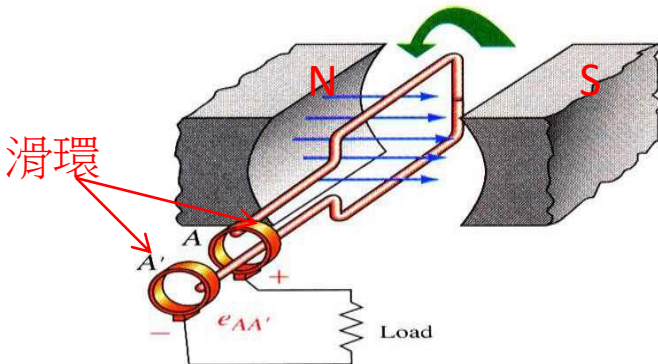
- 在靠近磁鐵凹槽處，裝上兩個繞在軟鐵芯上的線圈，分別對準磁鐵的兩極。搖轉線圈時，線圈導線中就感應出交流電。每當線圈轉半周，線圈對應的磁極就改變一次，電流方向也跟著改變。

- 發電機(generator)：一種將機械能轉換為電能的設備。基本上可分為直流發電機(DC generator)與交流發電機(AC generator)兩種，交流發電機又可分為單相(single phase)及三相(three phase)兩種。

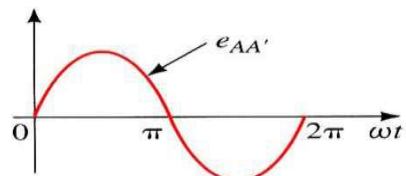
- 大型發電機：一般電力公司發電廠則有水力發電機(以水流推動水輪機轉動)、火力發電機(以燃燒油、煤或瓦斯方式產生高溫高壓蒸氣，推動蒸氣渦輪機轉動)、核能發電機(以原子分裂產生高溫高壓使蒸氣渦輪機轉動)等。發電廠的發電機通常為大容量(高達數仟MW)的三相交流同步發電機。

## 交流電的產生

· 交流電壓的產生：將一個矩形**線圈**放在磁鐵N極、S極所構成的固定**磁場**中，使該線圈以某一個特定方向、特定速度**旋轉**。線圈旋轉切割了磁通量 (magnetic flux) 使磁通量對時間發生變化，就發生了交變的感應電勢在導體中。



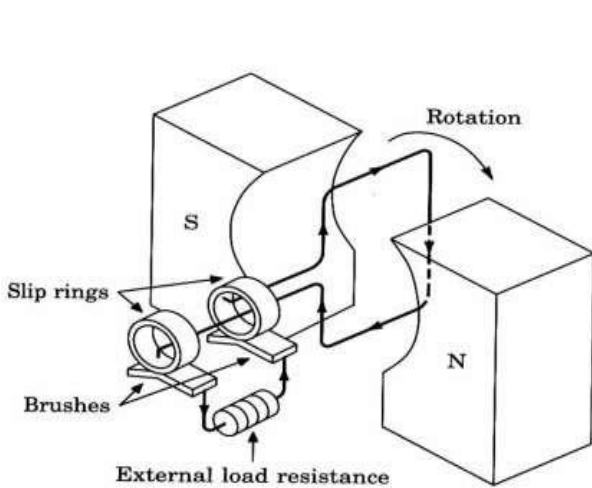
(a) Basic ac generator



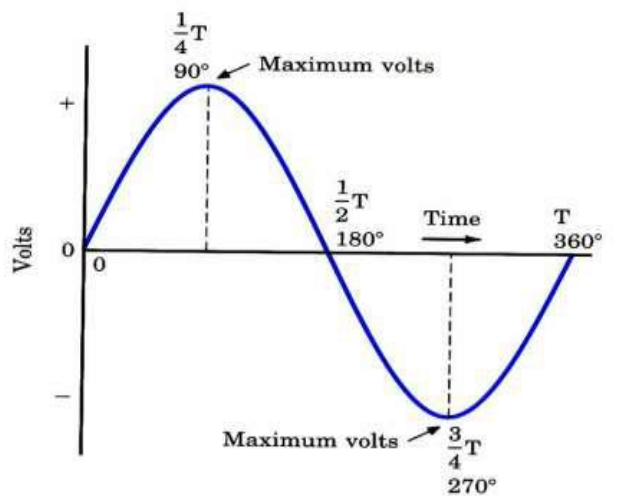
(b) Voltage waveform

## 單相交流發電機

· 線圈的兩個端點分別連接不同的**滑環**(slip-rings) ，該對滑環會隨線圈轉動而旋轉，將滑環以**電刷**(brushes)壓住，將電刷兩端加上負載(燈泡或電阻器)，則會有交變的電流通過負載，因此負載兩端會產生交變的電壓，此種電壓即為**單相**交流發電機的電壓。



(a)

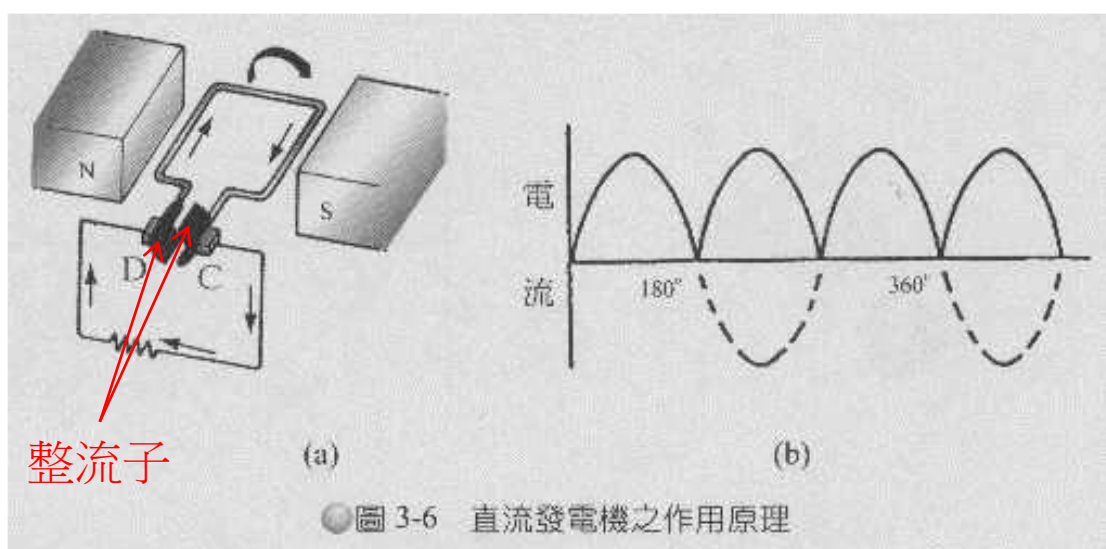


$$T = \text{Period} = \frac{1}{f} = \text{Time for 1 cycle}$$

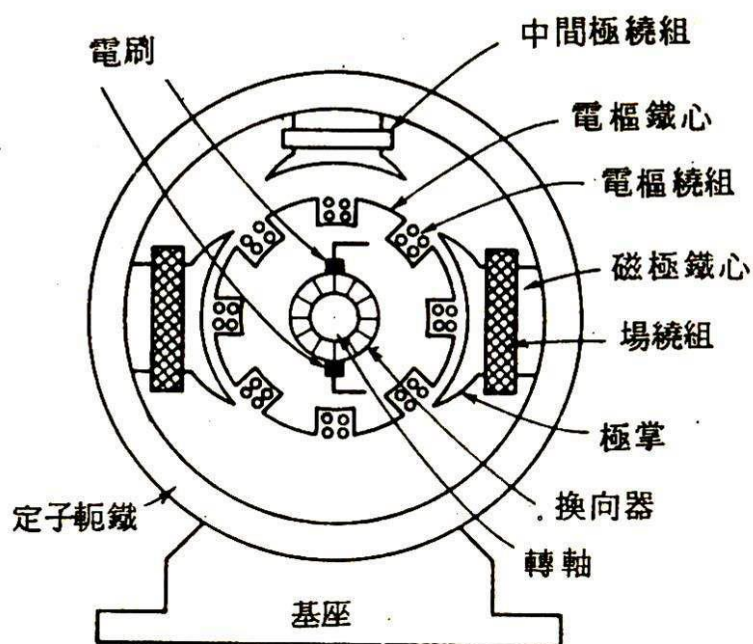
(b)

## 直流發電機的基本原理

· 直流發電機之作用原理，如圖 3-6所示，係在線圈兩端分別焊接在半圓形銅片上，此銅片俗稱**整流子**，再由**電刷**將電流由整流子引出。當線圈在磁場中轉動時，雖產生交流電，但整流子能使電流以一定的方向由C端電刷流出，D端電刷流入而成直流電



## 直流電機的結構圖



## The magnetic dipole

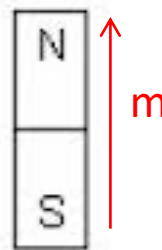
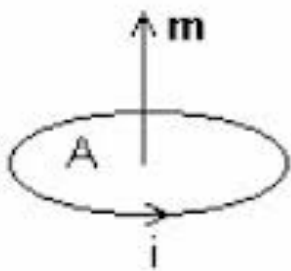
(What is the most elementary unit of magnetism)

A circular loop of a conductor carrying an electric current, which can generate a magnetic field.

A circular current loop can be considered the **most elementary unit of magnetism**.

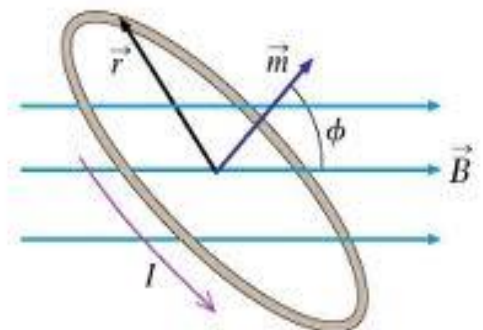
If a current loop has area  $A$  and carries a current  $i$ , then its **magnetic dipole moment** is  $m=iA$ .

The **unit** of magnetic moment:  $A \cdot m^2$  (amp  $\cdot$  meter<sup>2</sup>)



- The **torque** on a **magnetic dipole of moment  $m$**  in a magnetic **induction  $B$**  is then simply  $\vec{\tau} = \vec{m} \times \vec{B}$

- In free space  $\vec{\tau} = \mu_o \vec{m} \times \vec{H}$

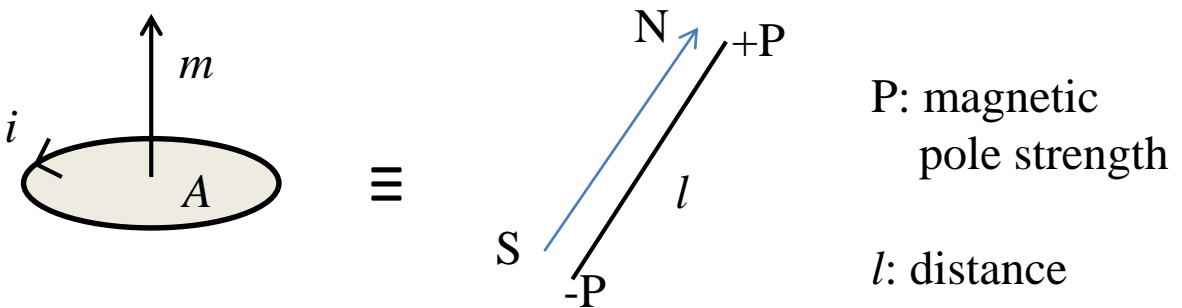


(-This mean that  $B$  tries to align the dipole so that the **moment  $m$**  lies **parallel to the induction**)

- The **energy** of the **dipole moment**  $\vec{m}$  in the presence of a magnetic induction (If no frictional forces are operating, the work done by the turning force will be conserved)  $E = -\vec{m} \cdot \vec{B}$

In free space  $E = -\mu_0 \vec{m} \cdot \vec{H}$

The field produced by a **current loop** is identical in form to the field produced by calculation from two hypothetical **magnetic poles** of strength separated by a distance  $l$ .



$$\vec{m} = p \cdot \vec{l}$$

- unit system in magnetism

The three unit systems { CGS - Gaussian  
 MKS - { Sommerfeld  
 Kennelly

Conversion factors:

- 1 oersted = 79.58 A/m
- 1 gauss =  $10^{-4}$  tesla
- 1 emu/cm<sup>3</sup> = 1000 A/m
- 1 maxwell =  $10^{-8}$  weber

- 1. The Earth's magnetic field is typically  $H=56$  A/m (0.70e)  
 $B=0.7 \times 10^{-4}$  tesla
- 2. The saturation magnetization of iron is  $M_0=1.7 \times 10^6$  A/m  
 $M_r=0.8 \times 10^6$  A/m
- 3. Lab electromagnet  $H = 1.6 \times 10^6$  A/m  
 $B = 2$  tesla

## 12 Magnetic fields

**Table 1.1** Principal unit systems currently used in magnetism

Quantity	SI (Sommerfeld)	SI (Kennelly)	EMU (Gaussian)
Field	$H$ A/m	A/m	oersteds
Induction	$B$ tesla	tesla	gauss
Magnetization	$M$ A/m	—	emu/cc
Intensity of magnetization	$I$ —	tesla	—
Flux	$\Phi$ weber	weber	maxwell
Moment	$m$ A m <sup>2</sup>	weber metre	emu
Pole strength	$p$ A m	weber	emu/cm
Field equation	$B = \mu_0(H + M)$	$B = \mu_0 H + I$	$B = H + 4\pi M$
Energy of moment (in free space)	$E = -\mu_0 m \cdot H$	$E = -m \cdot H$	$E = -m \cdot H$
Torque on moment (in free space)	$\tau = \mu_0 m \times H$	$\tau = m \times H$	$\tau = m \times H$

*Note:* The intensity of magnetization  $I$  used in the Kennelly system of units is merely an alternative measure of the magnetization  $M$ , in which tesla is used instead of A/m. Under all circumstances therefore  $I = \mu_0 M$ .

$H = 56$  A/m (0.7 Oe),  $B = 0.7 \times 10^{-4}$  tesla. The saturation magnetization of iron is  $M_0 = 1.7 \times 10^6$  A/m. Remanence of iron is typically  $0.8 \times 10^6$  A/m. The magnetic field generated by a large laboratory electromagnet is  $H = 1.6 \times 10^6$  A/m,  $B = 2$  tesla.

– Magnetic field calculation

– Magnetic field are usually produced by  $\left\{ \begin{array}{l} \text{solenoids} \\ \text{electromagnets} \end{array} \right.$

(solenoids are often cylindrical in shape)

$\left[ \begin{array}{l} \text{A solenoid is made by winding a large number of turns of insulated} \\ \text{Cu wire, or a similar electrical conductor, in a helical fashion on an} \\ \text{insulated tube know as a "former"}. \end{array} \right.$

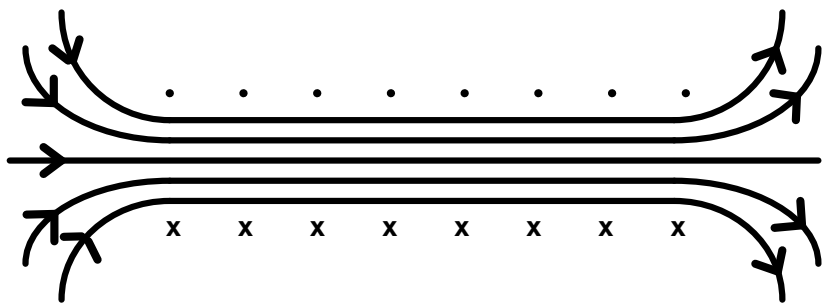
– The ferromagnetic core of an electromagnet generate a higher magnetic induction  $\vec{B}$  than a solenoid for the same magnetic field  $\vec{H}$

### 1. Field at the centre of a long thin solenoid

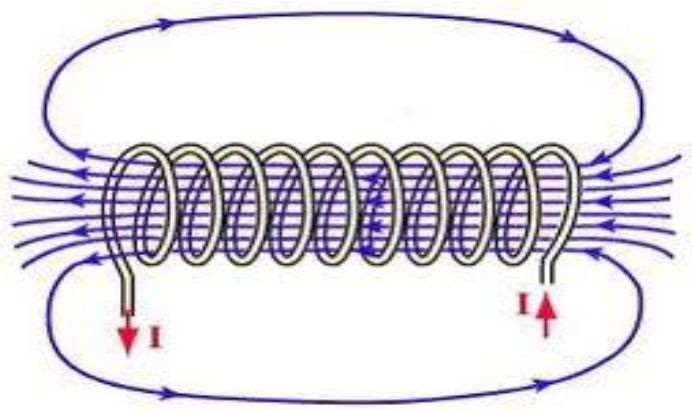
(what is the simplest way to produce a uniform magnetic field)

– If the solenoid has  $N$  turns wound on a former of length  $L$  and carries a current  $i$  amperes the field inside it will be

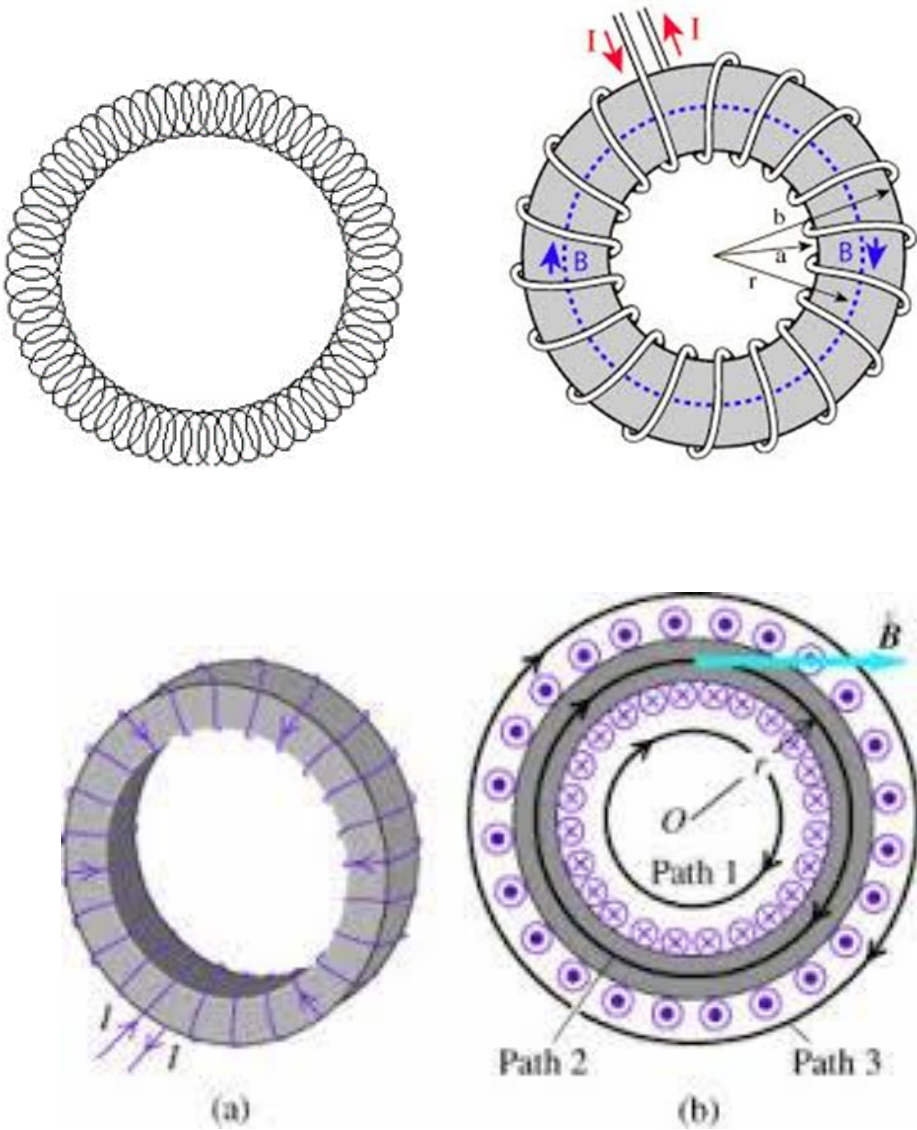
$$\vec{H} = \frac{Ni}{L} = ni \quad \left( n = \frac{N}{L} \right) \quad n: \text{the no. of turns per unit length}$$



Magnetic field lines Around a solenoid



A practical method of making an “infinite” solenoid is to make the solenoid toroidal in shape. This ensures that the field is uniform throughout the length of the solenoid.



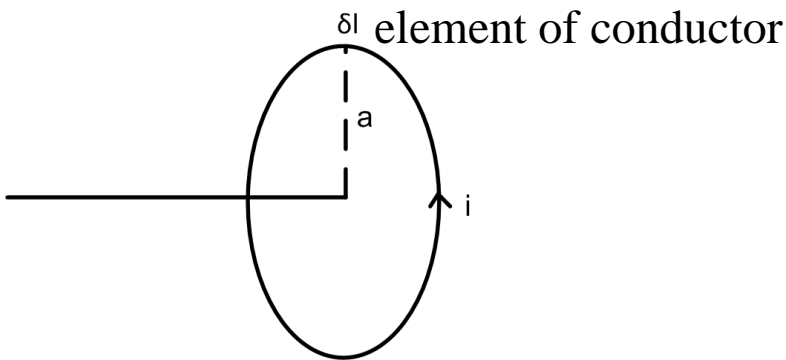
The magnetic field  $H = \frac{N}{2\pi r} i$   $\left\{ \begin{array}{l} N: \text{the total no. of turns} \\ r: \text{the radius of the toroid} \\ i: \text{the current flowing} \end{array} \right.$



## 2. Field due to a circular coil

(What is the field strength produced by the simplest form of coil geometry the single-turn)

(1) Field at the center of a circular coil



By Biot-Savart, 
$$\delta \vec{H} = \left( \frac{1}{4\pi r^2} \right) i \delta \vec{l} \times \vec{u}$$

$$H = \sum \frac{1}{4\pi r^2} i \delta l \sin \theta$$

$$\sum \delta l = 2\pi a, \quad d\vec{l} \perp \vec{u} \Rightarrow \theta = 90^\circ \Rightarrow \sin 90^\circ = 1$$

$$H = \frac{i}{2a} \text{ A/m}$$

## (2) Field on the axis of a circular coil

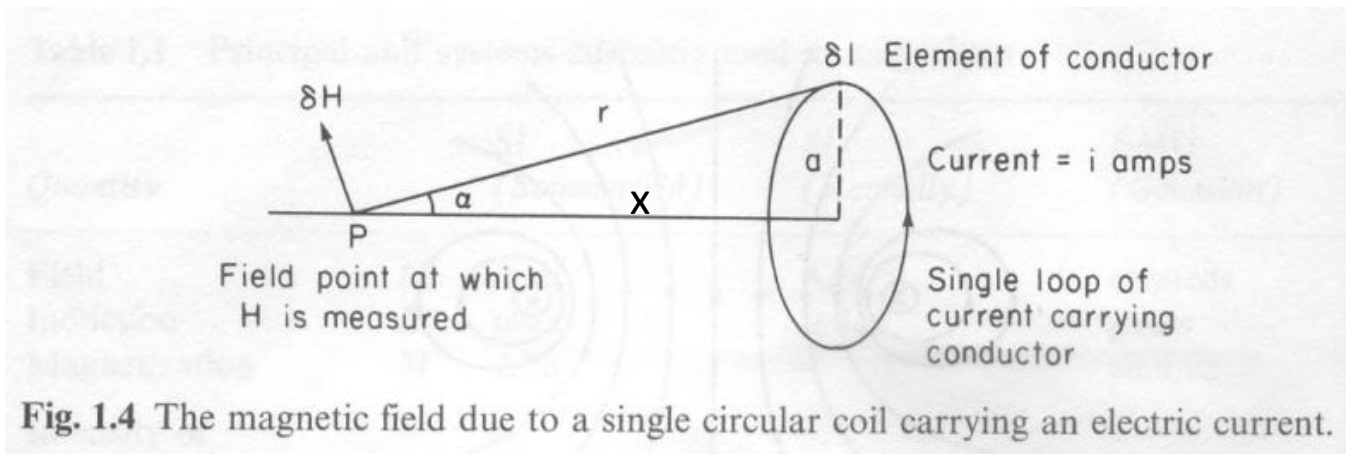
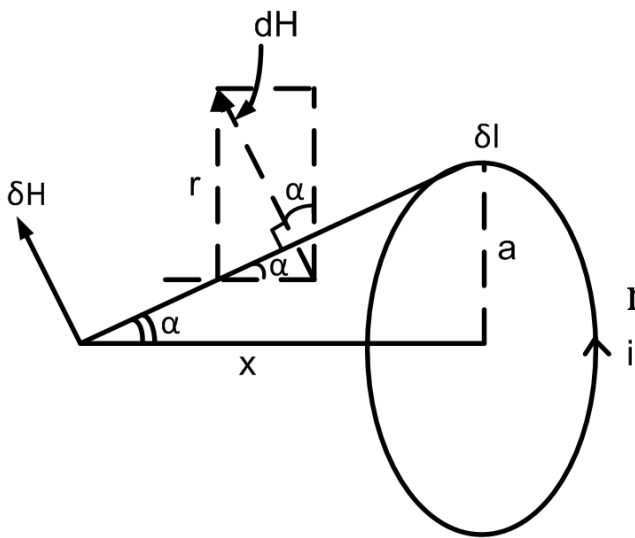


Fig. 1.4 The magnetic field due to a single circular coil carrying an electric current.



$$d\vec{H} = \frac{1}{4\pi r^2} i d\vec{l} \times \vec{u}$$

$\vec{u}$  is a vector long the  $r$  direction

$$r = \frac{a}{\sin \alpha}, \quad d\vec{H} = \frac{1}{4\pi a^2} (\sin^2 \alpha) i d\vec{l} \times \vec{u}$$

$$d\vec{H}_{axial} = d\vec{H} \sin \alpha$$

$$d\vec{H}_{axial} = \frac{1}{4\pi a^2} (\sin^3 \alpha) i d\vec{l} \times \vec{u}$$

$$\int dl = 2\pi a, \quad d\vec{l} \perp \vec{u}$$

$$\vec{H}_{axial} = \frac{i}{2a} \sin^3 \alpha = \frac{i}{2a} \left(\frac{a^3}{r^3}\right) = \frac{ia^2}{2r^3} \quad r = \sqrt{a^2 + x^2}$$

$$\vec{H}_{axial} = \frac{ia^2}{2(a^2 + x^2)^{3/2}}$$

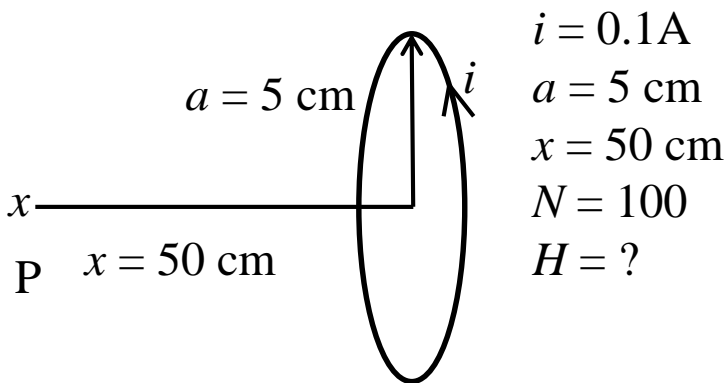
This can be expressed in the form of a series in  $x$  (and by symmetry all terms of odd order must have zero coefficients so the form of the dipole field become)

$$H = H_o (1 + C_2 x^2 + C_4 x^4 + C_6 x^6 + \dots)$$

$$H_o = \left(\frac{i}{2a}\right) \text{ (the field at the center of the coil)}$$

$$C_2 = -\frac{3}{2a^2}, \quad C_4 = -\frac{15}{8a^4}, \quad C_6 = -\frac{105}{48a^6}$$

Ex 1.3 If a coil of 100 turns and diameter 10 cm carries a current of 0.1A, calculate the magnetic field at a distance of 50 cm along the axis of the coil.



$i = 0.1 \text{ A}$   
 $a = 5 \text{ cm}$   
 $x = 50 \text{ cm}$   
 $N = 100$   
 $H = ?$

$$\begin{aligned}
 H &= \frac{100(0.05)^2(0.1)}{2\left[(0.05)^2 + (0.5)^2\right]^{3/2}} \\
 &= \frac{0.025}{2(0.253)^{3/2}} \\
 &= 0.0098 \text{ A/m}
 \end{aligned}$$

(3) off-axis field a circular coil  $\left( \begin{array}{l} \text{In the vast majority of cases there is no} \\ \text{closed-form analytic solution for the field} \\ \text{generated by a current-carrying conductor} \end{array} \right)$

In the case of the off-axis field of the single circular loop the analysis leads to an elliptic integral which has no exact solution.

By Biot - Savart law  $dH = \frac{idl \times u}{4\pi r^2}$   $r$  : the distance from the coil

$$dH = \frac{idl \sin \theta}{4\pi(x^2 + a^2)}$$

$\left( \begin{array}{l} \text{where now } a \text{ is also a function of } \theta \text{ instead of being a constant. In the case of} \\ \text{the off-axis field, the field strength can be calculated from this eq. by a computer} \\ \text{using numerical techniques.} \end{array} \right)$

## Chap 2 Magnetization and Magnetic Moment

(When going on to consider magnetic materials it is first necessary to define quantities which represent the response of these materials to the field. These quantities are **magnetic moment** and **magnetization**. Once that has been done, we can consider another property. The susceptibility, which is closely related to the permeability.)

- Consider the effect that a magnetic materials has on the magnetic induction  $\underline{B}$  when a field pass throught it.

↓

Magnetization

- Materials can alter the magnetic induction ( $B$ )
  - larger → (paramagnets, ferromagnets)
  - Smaller → (diamagnet)

- Magnetic moment ( $m$ )

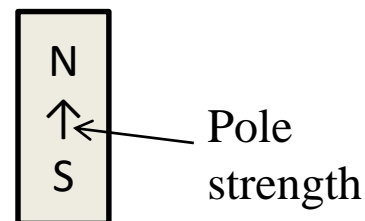
(can we use the torque on a specimen in a field of know strength to define its magnetic properties?)

- The torque on the dipole in the presence of a magnetic field in free space is given by  $\vec{\tau} = \vec{m} \times \vec{B}$

- The  $\vec{m}$  can be expressed as the maximum torque on a magnetic dipole  $\tau_{\max}$  divided by  $B$

$$\vec{m} = \frac{\vec{\tau}_{\max}}{\vec{B}}$$

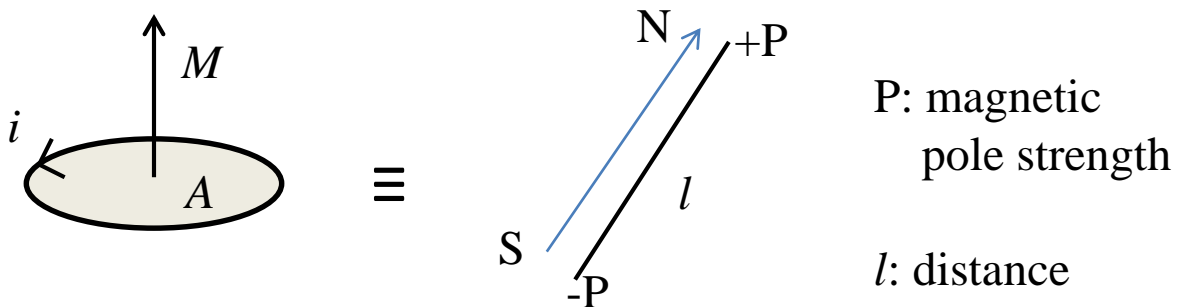
$$= \frac{\vec{\tau}_{\max}}{\mu_0 \vec{H}} \quad (\text{in free space})$$



- Unit of  $m$ :  $A \cdot m^2$  (ampere·meter<sup>2</sup>)
- This formula applies equally to a current loop or a bar magnet

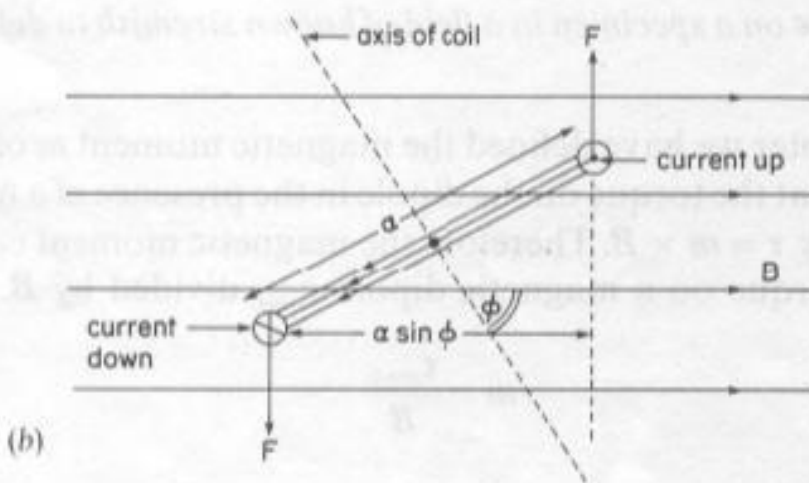
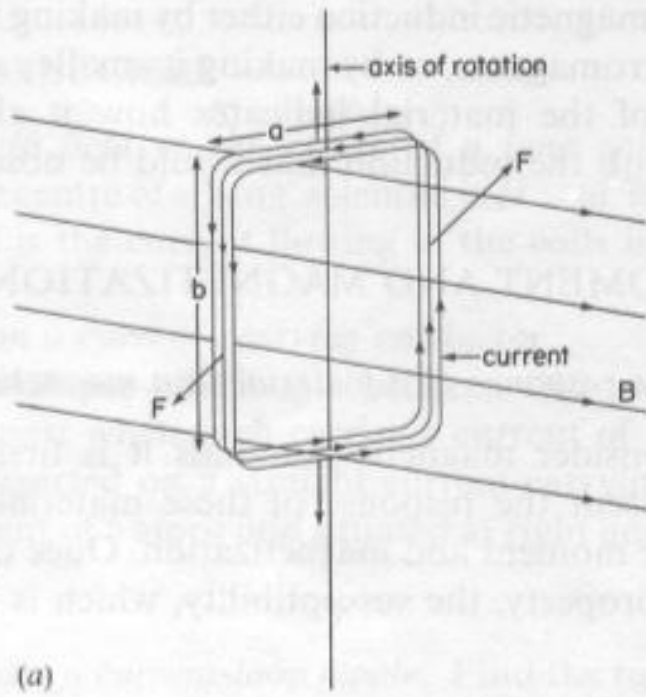
## The unit of magnetic moment

- The magnetic moment of  $1 \text{ ampere}\cdot\text{meter}^2$  experiences a maximum torque of  $1 \text{ newton}\cdot\text{meter}$  when oriented perpendicular to magnetic induction of  $1 \text{ tesla}$ .

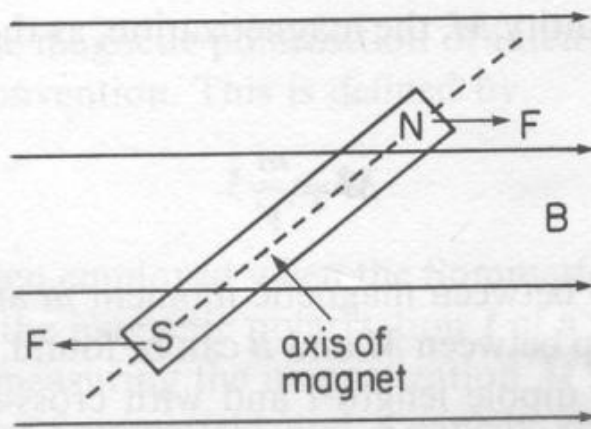


In the case of current loop       $m=iA$        $i$ : the current flowing (A)  
 $A$ : the cross-sectional area of the loop ( $\text{m}^2$ )

In the case of a bar magnet       $m=pl$        $p$ : the pole strength ( $\text{A}\cdot\text{m}$ )  
 $l$ : the dipole length (m)

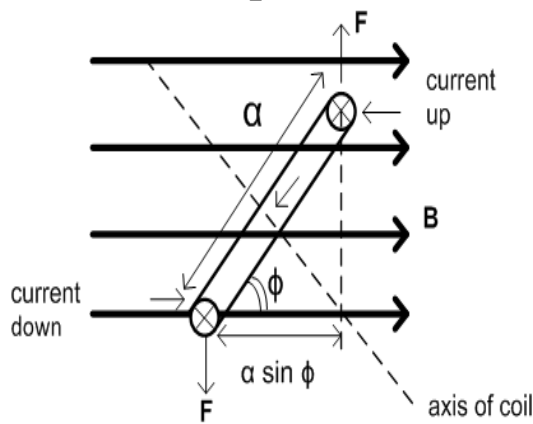


**Fig. 2.1** The torque on a current loop in an external magnetic field; (a) side view, and (b) to view. If the loop is free to rotate the torque turns the loop until its plane is normal to the field direction.



**Fig. 2.2** The torque on a bar magnet in an external magnetic field. If the bar is free to rotate the torque turns the bar until its plane is parallel to the field direction.

In a current loop



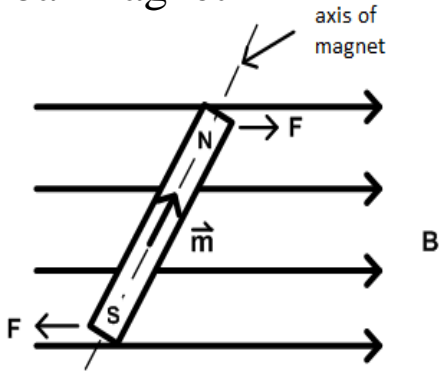
$$\therefore E = -\vec{m} \cdot \vec{B}$$

$\therefore$  當  $\vec{m} // \vec{B}$  時， $E_{\min}$  會產生

The torque turns the loop until its plane is normal to the field direction.

$$m = iA$$

In a bar magnet



The magnetic moment vector  $\vec{m}$  in a bar magnet tends to align itself with  $B$  under the action of the torque.

$$\therefore E = -\vec{m} \cdot \vec{B}$$

$\therefore$  當  $\vec{m} // \vec{B}$  時， $E_{\min}$  會產生

A torque aligns the magnet parallel to the local direction  $B$ .

$$m = p \cdot l$$

$p$ : amp·meter

$l$ : dipole length

$$M = \frac{m}{v}$$

↖ Magnetic moment  
↖ Volume of a solid

In the Sommerfeld convention

this is given

$$p = \frac{\Phi}{\mu_0}$$

$$m = pl$$

$$= \frac{\Phi l}{\mu_0}$$

$\Phi$  is the flux in webers passing through the current loop or bar dipole.  $l$  is the dipole length.

- The “pole strength” is arising from the more traditional CGS Treatment of magnetism, in which pole strength was defined in terms of the magnetic flux  $\Phi$  emanating from a single magnetic pole



## Magnetization ( $M$ )

–How are the magnetic properties of the material and the magnetic induction  $B$  related?

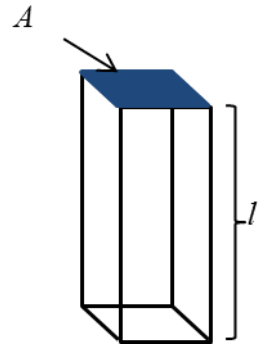
$$M, \text{ magnetization,} = \frac{\text{the magnetic moment}}{\text{per unit volume}} \quad M = \frac{m}{V}$$

–A bar magnet with flux density  $\phi$  at the center, dipole length  $l$  and with cross-sectional area  $A$  has a magnetic moment  $m$ , given by  $m = \frac{\phi l}{\mu_0}$

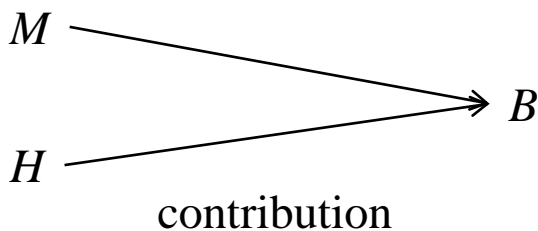
–The magnetization  $M$  is given by

$$M = \frac{m}{Al} = \frac{\Phi}{\mu_0 A} \quad (m = \frac{\Phi l}{\mu_0})$$

$$= \frac{B}{\mu_0} \quad (\Phi = B \cdot A)$$



– In this case there are **no conventional external electric currents** present to generate an **external magnetic field ( $H$ )**, so  $B = \mu_0 M$



–If both  $M$  and  $H$  are present then their contributions can be summed.

## Relation between $H$ , $M$ , $B$

(Can we define a **universal equation** relating these three magnetic quantities  $H$ ,  $M$ ,  $B$ ?)

– Magnetic induction  $B$  consists of two contribution  $\left\{ \begin{array}{l} H \\ M \end{array} \right.$   

$$B = \mu_0(H + M)$$

$\left\{ \begin{array}{l} \text{In free space } B = \mu_0 H \text{ (The magnetic induction in free space)} \\ \text{In a material } B = \mu_0 M \text{ (The contribution to the induction from the magnetization of a material)} \end{array} \right.$

Unit:  $B$  – tesla

$$H - \text{A/m}$$

$$M - \text{A/m}$$

–  $H$  is generated from  $\left\{ \begin{array}{l} \text{solenoid} \\ \text{electromagnet} \\ \text{permanent magnet} \end{array} \right.$  outside the material  
 (by electrical currents)

–  $M$  is generated by  $\left\{ \begin{array}{l} \text{spin} \\ \text{orbital angular momentum} \end{array} \right.$  of electrons within the solid

– A related quantity, the magnetic polarization or intensity of magnetization  $I$  is used in the Kennelly convention. This is defined by  $I = \mu_0 M$

unit of  $I$  : tesla